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GENERALIZED GENERAL RELATIVISTIC QUANTUM STATIC FIELD EQUATION AND QUANTIZATION OF ENERGY EQUATION

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ABSTRACT

The Hamiltonian and the momentum for full spherical coordinates are related to each of her. Using conventional particle wave duality, a relativistic quantum equation based on energy momentum tensor of generalized general relativity was derived. The solution shows that the energy and angular momentum are quantized. It also predicts graviton and gravity wave existence.

INTRODUCTION THE HISTORY OF COSMOLOGY

The history of cosmology starts from the beginning of human life on The earth. People observe the sun moving and distributing light from the morning. At night the moon appear on the sky as a golden hemisphere surrounded by glittery beautiful objects. Some of us ask; what is the nature of these astronomical objects? Why some of them move and look brighter than others [1, 2, 3].The accumulated knowledge for these objects led scientists like Kepler to formulate some rules for some of the regular behaviors of these astronomical objects these rules are concerned with the motion of planets around the sun [4, 5]. Later on lsaag Newton discovered that is called, the gravitational field, is responsible for the motion planets around the sun [6].Gravitational field is used to explain the influence of a massive body that spreads in to the space around it the so called a field, which produces a force on another massive body. Very close. [7, 8].Newton Lows of gravitational succeeded in describing the motion of macroscopic objects [9] until the beginning of twentieth century, where Michelson-Morley experiment indicates the violation of the Newton Low of addition of velocity for light [10, 11, and 12].This experiment shows that the speed of light in vacuum is always constant and is completely independent on the motion of the source or observer [13].

Unfortunately Newton gravitational Low suffers from noticeable setbacks. For instance it fails to describe the prehclian of mercury, beside the failure in describing the behavior of quasi-stellar objects and black holes [14, 15].

In 1915, Albert Einstein developed his theory of general relativity, having earlier shown that gravity dos influence light's motion [16, 17]. Einstein's theory of general relativity (GR) is one of the fundamental physical theories at the present time; it describes a number of gravitational physical phenomena, which agree with astronomical observations [18, 19].Despite these successes GR suffers from being isolated from the main stream of physics. This is since the equation of motion and the energy momentum tensor conservation Lows differs radically from that in other physical theories, which are derived from the action principle. It also suffers from the lack of a full expression for the energy-momentum tensor of the gravitational field [20, 21, 22,].General Relativity Lows can't easily explain the behavior of exotic astronomical objects, like black holes, pulsars, quasars and neutron stars [23, 24, and 25].For instance it is difficult to explain the large red shift of quasars within the frame work of GR [26, 27, and 28].The behavior of black holes is even more complex. A black holes is a place where gravity is very

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strong that even light can't get out. The gravity is so strong because matter has been squeezed in to a tiny space this can happen when a star is dying. Because not light can get out, people can't see black holes. They are invisible, space telescopes with special tools help find black holes. The special tools can see how stars that are very close to black holes act differently than other stars [29, 30].Black holes can be big or small. Scion tests think the smallest black holes are as large mountain. Another kind of black holes is called staller its mass can be up to 20 times than the mass of the sun. There may be many, many staller mass black holes in Earth's galaxy. Stiller black holes are made when the center of Avery big star falls in upon itself, o collapse. When this happens, it causes a super move. A super move is an exploding star that blasts part of the star in to space. All these gravitational pheromone seems to reed a full quantum gravitational theory as pointed out by many anthers [29,30].Einstein general relativity GR is isolated in its geometrical content from the main stream of physics. This is because there is no well-established quantum gravitational theory. Thus the behavior of black holes and neutron stars can't be explained fully. More over the unification of gravity with other field is not yet achieved. Different attempts were made to find a full quantum gravitational theories [29, 30, 31, 32, 33, 34, 35].In some of them wheeler de Witt quantum wave function is used to describe the universe evolution [36, 37, 38, 39, 40].In to the approaches canonical quantization based on GR is also proposed [41, 42].A quantum model based on quantized general relativity (GGR) is also proposed by some people [43]. The quantum models based on GR does not based on the Hamiltonian which is not defined in GR. The ones based on GGR do not have a wave function and do not dens be isolated bodies. Motivated by the successes of quantum seismological models [44], the aim of this work is devoted to find quantum mode to quantized static field generated by isolated stars. This can describe the behavior of black holes and quasi-stellar and to remove singularity.

THE FULL SPHERICAL QUANTUM GRAVITY EQUATION

In the system of units where $c \neq 1$ the (quantum general equation) becomes $\hat{E}\psi = c\hat{P}\psi + f(r)\psi$ (1)

This equation can be rewritten by taking in to account the fact that in (classical mechanic), the energy is given by:

$$
E = \int F \cdot dr = \int m \frac{dv}{dt} \cdot dr = m \int d\underline{v} \cdot \frac{dr}{dt}
$$

$$
E = m \int \underline{v} \cdot d\underline{v} = \frac{mv^2}{2}
$$
(2)

But if the system is oscillating the velocity can thus give by $v(t) = v_m sin\omega t$ (3) Where the effective velocity is given by:

$$
v = \frac{v_{max}}{\sqrt{2}} = \frac{v_m}{\sqrt{2}}, \quad V^2 = \frac{v_m^2}{2}
$$
(4)
If one rewrite (5.5.2) to be
\n
$$
E = \frac{mv_m^2}{2}
$$
(5)
\nIt follows that
\n
$$
E = mv^2
$$
(6)
\nAlternatively for harmonic oscillator
\n
$$
E = \frac{1}{2} m\omega^2 A^2 = \frac{1}{2} m v_m^2 = m(\frac{V_m}{\sqrt{2}})^2 = m v_e^2 = m v^2
$$
(7)
\nThus
\n
$$
E = mv^2 = (mv) \cdot v = P \cdot v
$$
(8)
\n
$$
E = P_x V_x + P_y V_y + P_z V_z
$$
(9)
\nFor light
\n
$$
V_x = V_y = V_z = c
$$
(10)
\nThus
\n
$$
E = cP_x + cP_y + cP_z
$$
(11)
\nIn spherical coordinate
\n
$$
E = cP_r + cP_\theta + cP_\phi
$$
(12)
\nIn view of (12) equation (1) can be written as
\n
$$
E = cP_r + cP_\theta + cP_\phi + f(r)
$$
(13)
\nMultiply both sides by ψ to E get

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 $E\psi = cP_r\psi + cP_\theta\psi + cP_\phi\psi + f(r)\psi$ (14) Replacing E and P by their corresponding operates one gets $\hat{E}\psi = c\hat{P}_r\psi + c\hat{P}_\theta\psi + c\hat{P}_\phi\psi + f(r)$ (15) But the energy and momentum operators take the form $\hat{E} = i \frac{\partial}{\partial \theta}$ ∂t $\hat{P}_r = \frac{\hbar}{i}$ i ∂ дţ $\hat{P}_{\theta} = -\frac{\hbar}{i\omega}$ ir $\frac{\partial}{\partial \theta} \hat{P}_{\phi} = \frac{\hbar}{i}$ i 1 $rsin\theta$ ∂ $\frac{1}{\partial \phi}$ (16) Thus the full quantum equation becomes $i\hbar \frac{\partial \psi}{\partial t} = \frac{\hbar}{i}$ $\frac{\hbar}{i}\left[c\frac{\partial \psi}{\partial r} + \frac{c}{r}\right]$ r $\frac{\partial \psi}{\partial \theta} + \frac{c}{r \sin \theta}$ $rsin\theta$ $\frac{\partial \psi}{\partial \phi} + f(r)\psi$ $i\hbar \frac{\partial \psi}{\partial t} = \frac{\hbar}{i}$ $\frac{\hbar}{i}\left[c\frac{\partial \psi}{\partial r} + \frac{c}{r}\right]$ r $\frac{\partial \psi}{\partial \theta} + \frac{c}{r \sin \theta}$ $rsin\theta$ $\frac{\partial \psi}{\partial \phi} + f(r)\psi$ (17) Using separation of variables $\psi(r, \theta, \phi, t) = \tilde{R}(r) + \theta \Phi(\phi) v(t)$ (18) Equation (17) becomes $iR\Phi\Theta\frac{dv}{dt} = \frac{\hbar c}{i \times}$ $\frac{\hbar c}{i \times i} v H \Theta \Phi \frac{\partial \ddot{R}}{\partial r}$ $\frac{\partial \ddot{R}}{\partial r} + \frac{\hbar c}{i \times}$ $\frac{\hbar c}{i \times i} v \tilde{R} \Phi \frac{\partial \theta}{\partial \theta} + \frac{c \hbar}{ir sin \theta} \tilde{R} \theta \frac{\partial \Phi}{\partial \phi} + f(r) R \theta \Phi v$ Divide by $R \theta \phi \nu$ iħ \mathcal{V} $\frac{dv}{dt} = \frac{c\hbar}{i}$ i 1 ̃ ∂R $\frac{\partial \tilde{R}}{\partial r} + \frac{c\hbar}{i}$ i 1 $r\theta$ $\frac{\partial \theta}{\partial \theta} + f(r) + \frac{c\hbar}{i}$ i 1 $r\phi$ sin θ $\frac{\partial \Phi}{\partial \phi} = C_0 = E$ Thus iħ ∂v $\frac{\partial}{\partial t} = E$ (19) $f(r) + \frac{c\hbar}{t}$ i 1 ̃ $d\tilde{R}$ $\frac{d\tilde{R}}{\partial r} + \frac{ch}{i}$ i 1 $r\theta$ $\frac{\partial \theta}{\partial \theta} + \frac{ch}{i\phi (crs)}$ $i\phi(crsin\theta)$ $\frac{\partial \Phi}{\partial \phi} = E$ Multiply by cħ i $d\tilde{R}$ $\frac{d\tilde{R}}{dr} + f(r)r + \frac{c\hbar}{i\Theta}$ $i\theta$ $\frac{d\theta}{d\theta} + \frac{ch}{\phi(isi)}$ $\phi(\textit{isin}\theta)$ $\frac{d\Phi}{d\phi} = Er$ cħ i r $\frac{r}{R}$ + $(f(r) - E)r = \frac{ch}{i}$ i $\frac{d\theta}{d\theta} + \frac{ch}{\phi}$ ɸ 1 $sin\theta$ $\frac{d\Phi}{d\phi} = C_1$ Thus ℎ i r ̃ $d\tilde{R}$ $\frac{d}{dr} + [f(r) - E]r = C_1$ (20) cħ iΘ $\frac{d\theta}{d\theta} + \frac{ch}{\Phi i}$ ɸ 1 sinθ $rac{d\Phi}{d\phi} = C_1$ (21) Multiply by $(sin\theta)$ cħ i sinθ θ $\frac{d\theta}{d\theta} + \frac{ch}{\phi}$ ɸ $\frac{d\Phi}{d\phi} = C_1 sin\theta$ (22) $c\hbar \frac{\sin\theta}{\cos\theta}$ iΘ $\frac{d\theta}{d\theta} - C_1 sin\theta = -\frac{ch}{\phi i}$ ɸ $\frac{d\Phi}{d\phi} = C_1$ (23) Thus cħ i sinθ θ $\frac{d\theta}{d\theta} - C_1 sin\theta = C_2$ (24) The function $f(r)$ can be found from equations (3),(4) where the GGR Hamiltonian is given by $H = \alpha R^2 + \alpha \frac{\dot{B}}{AB} \dot{R} = \frac{1}{3}$ $\frac{1}{3} \alpha R^2 + \frac{1}{3}$ $\frac{1}{3} \alpha R^2 + \frac{1}{3}$ $\frac{1}{3}\alpha R^2 + \alpha \frac{\dot{B}}{AB} \dot{R}$ (25) The corresponding momentum components in spherical coordinates is given by 1 $\frac{1}{3}P_r = \frac{1}{3}$ $\frac{1}{3}\alpha R^2 + \frac{2}{3}$ 3 $\alpha\ddot{R}$ $\frac{\alpha \ddot{R}}{A} - \frac{\alpha}{3}$ 3 ÀŘ R^2

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1 $\frac{1}{3}P_{\theta} = \frac{1}{3}$ $rac{1}{3}\alpha R^2 + \frac{\dot{R}}{3r}$ $3rA$ 1 $\frac{1}{3}P_{\phi} = \frac{1}{3}$ $\frac{1}{3}\alpha R^2 + \frac{\dot{R}}{3r}$ $\frac{1}{3rA}$ (26) Thus inserting equation (26) in equation (25) yields $H = \frac{1}{2}$ $\frac{1}{3}[P_r + P_\theta + P_\phi] - \frac{2}{3}$ 3 $\alpha\ddot{R}$ $\frac{\alpha \ddot{R}}{A} + \frac{\alpha}{3}$ 3 $\dot A$ $\frac{\dot{A}}{A^2}\dot{R}-\frac{2}{3}$ 3 Ŕ $\frac{1}{rA}$ + $\frac{\alpha \dot{B}}{AB} \dot{R}$ (27) In the system of units where $c \neq 1$ The Hamiltonian becomes $H=\frac{c}{2}$ $\frac{1}{3}[P_r + P_\theta + P_\phi] + f(r)$ (28) Or by following Dirac relativistic quantum equation approach $H = c\alpha_r P_r + \alpha_\theta P_\theta + \alpha_\phi P_\phi + f(r)$ (29) Where $2 \alpha \ddot{R}$ $\alpha \dot{A}$ $2 \dot{R}$ $\alpha \dot{B}$

$$
f(r) = -\frac{2}{3}\frac{\alpha R}{A} + \frac{\alpha A}{3}\frac{A}{A^2}\dot{R} - \frac{2}{3}\frac{R}{rA} + \frac{\alpha B}{AB}\dot{R}
$$

$$
f(r) = -\frac{2}{3}\frac{\alpha R}{A} + \frac{\alpha A}{3}\frac{\dot{A}}{A^2}\dot{R} - \frac{2}{3}\frac{R}{rA} + \frac{\alpha B}{AB}\dot{R}
$$
(30)

SOLUTION OF THE RADIAL PART

Since most of astronomical objects have spherical shape. Therefore it is suitable to use spherical coordinates. In view of equation (20) the radial part is given by

$$
\frac{c\hbar}{i R} \frac{r}{R} \frac{d\bar{R}}{dr} + [f(r) - E]r = C_1
$$
\n(31)
\nThus, separation of R and r dependent parts yield
\n
$$
\int \frac{d\bar{R}}{\bar{R}} = \frac{iC_1}{c\hbar} \int \frac{dr}{r} + \int \frac{i}{c\hbar} (E - f(r)) dr + C_4
$$
\n
$$
\ln \bar{R} = \frac{iC_1}{c\hbar} \ln r + \frac{i}{c\hbar} [Er - \int f(r) dr] + C_4 \ln \bar{R} - \ln r^{ic_3} = \frac{i}{c\hbar} [Er - \int f(r) dr] + C_4
$$
\n
$$
\bar{R} = C_5 r^{ic_3} e^{\frac{i}{c\hbar} [Br - f(r) dr]} \qquad (32)
$$
\nWhere
\n
$$
C_3 = \frac{C_1}{c\hbar} C_5 = e^{c_4}
$$
\n(33)
\nThe radial wave function can be rewritten in the form
\n
$$
\bar{R} = e^{i c_3 \ln r} e^{\frac{i}{\hbar} [Br - f(r) dr]} \qquad (34)
$$
\n
$$
R = R_0 = constant \qquad (35)
$$
\nWhere
\n
$$
H = \alpha R_0^2 = constant \qquad (36)
$$
\nBut since the energy density is equation on energy multiplied by the number of them, therefore
\n
$$
H = \hbar \omega |\psi|^2 = \hbar \omega |\bar{R}|^2 = constant \qquad (37)
$$
\n
$$
T = \hbar \omega |\psi|^2 = \hbar \omega |\bar{R}|^2 = constant \qquad (37)
$$
\n
$$
T = \hbar \omega |\psi|^2 = \hbar \omega |\bar{R}|^2 = constant \qquad (38)
$$
\nFor simplicity let
\n
$$
C_1 = 0 \qquad C_3 = 0 \qquad (39)
$$
\nFor simplicity let
\n
$$
C_1 = 0 \qquad C_3 = 0 \qquad (40)
$$
\nIn view of equation (4) and (8)
\n
$$
\frac{i}{e^{\pi E r_0}} = c_7 \qquad (41)
$$
\nBut, since the number of gravitons are constant, and from (6) which hews that R is real as for as the Hamiltonian
\n(energy) is real thus

 $R = \cos \frac{Er_0}{r}$ $\frac{Er_0}{c\hbar}$ + isin $\frac{Er_0}{c\hbar}$ $\frac{Er_0}{c\hbar}=c_7$, $\cos\frac{Er_0}{c\hbar}$ $\frac{Er_0}{c\hbar} = c_7 \sin \frac{Er_0}{c\hbar}$ $\frac{\partial}{\partial t} = 0$ (42) Hence

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 Er_0 $\overline{c}\hbar$ $= 2n\pi$ $E =$ $2n\pi c\hbar$ r_0 (43) $E=$ nch r_0 (44) According to Bohr model the minimum Bohr radian corresponds to $2\pi r_0 = \lambda$ Thus $E =$ $nc\hbar(2\pi)$ λ $E = 2\pi n f$ (45)

SOLUTION OF CONSTANT SCALAR CURVATURE

There are many useful solutions which can be found when the mass density is constant. According to GR it implies constant R using equation (30)

$$
f(r) = -\frac{2}{3}\alpha \frac{\ddot{R}}{A} + \frac{\alpha \dot{A}}{3A^2}\dot{R} + \frac{\alpha \dot{B}}{AB}\dot{R} - \frac{2}{3}\frac{\dot{R}}{rA}
$$
 (46)
For constant *R*
 $R = R_0$ $\dot{R} = 0$ $\ddot{R} = 0$ (47)
Thus (1) given (2)
 $f(r) = Zero$
Substituting in (17) and consider ψ to depend on *r* only for afield generated by a spherical body.
 $\psi = \psi(r)$ (48)
 $\frac{\partial \psi}{\partial \theta} = Zero$ (49)
Then
 $i\hbar = \frac{\partial \psi}{\partial t} = \frac{c\hbar}{i} \frac{\partial \psi}{\partial r} + Zero + Zero$
 $\therefore \frac{\partial \psi}{\partial t} = \frac{c\hbar}{i} \frac{\partial \psi}{\partial r} + Zero$ (52)

$$
i\hbar = \frac{V}{\partial t} = \frac{V}{i \partial r}
$$
 (50)
Substituting

$$
\psi = A \sin(kr - \omega t)
$$
 (51)

$$
\frac{\partial \psi}{\partial t} = -\omega A \cos(kr - \omega t) \frac{\partial \psi}{\partial r} = kA \cos((kr - \omega t) (53))
$$

Substituting in
\n
$$
i\hbar \frac{\partial \psi}{\partial t} = \frac{hc}{i} \frac{\partial \psi}{\partial r}
$$
\n
$$
-i\hbar \omega A \cos((kr - \omega t), \frac{hc}{i} kA \cos((kr - \omega t))
$$
\n(52)

 $-i\hbar\omega = c k \frac{\hbar}{\hbar}$ i Multiply by \hat{i} and divide by \hbar $W = CK$ (53) This indicates that the solution is consistent with the ordinary relation between angular frequency and wave number. Thus $\psi = A \sin((kr - \omega t)$ (54) Is a solution which indicates that gravitons are travelling waves moving with speed of light. Using the relation $cos^{2} + sin^{2} = 1$ The probability of existence of particles is $|\psi|^2 = A^2 \cos^2 \theta$ (55) But $\cos 2\theta = \cos^2 \theta - \sin^2 \theta = -2\cos^2 \theta - 1$ $\cos^2\theta = \frac{1}{2}$ $\frac{1}{2}$ [cos2 θ + 1]

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 $h = |\psi|^2 \frac{A^2}{2}$ $\frac{1}{2}[1 + cos2\theta(kr - \omega t)]$ (56) This means that a gravitational wave can be propagated with wave function $\psi = A\cos(kr - \omega t)$ (57) The intensity of waves is $n = \frac{A^2}{2}$ $\frac{1}{2}[1 + \cos 2(kr - \omega)]$ $n = n_0 + \frac{A^2}{2}$ $\frac{1}{2}\cos 2(kr - \omega t)$ (58) Where $n_0 = \frac{1}{2}$ A^2 (59) The constant form n_0 can be considered as. A back ground constituting vacuum energy.

GRAVITATIONAL FIELD QUANTIZATION

The gravitational field for spherically symmetric body satisfies the equation $i\hbar \frac{\partial \psi}{\partial t} = \frac{\hbar c}{i}$ i $\partial \psi$ $\frac{\partial \mathbf{r}}{\partial r}$ (60) Separating the variables in to time and radial function $\psi = uv = u(t)v(r)$ (61) Thus $i\hbar \frac{v\partial u}{\partial t} = \frac{\hbar c}{i}$ i $\frac{u\partial v}{\partial r}$ (62) Dividing both sides by uv yield iħ \overline{u} $\frac{\partial u}{\partial t} = \frac{\hbar c}{i}$ i 1 \mathcal{V} $\frac{\partial v}{\partial r} = C_0 = E$ (63) Thus $i\hbar \frac{\partial u}{\partial t} = Eu$ (64) Which is the energy Eigen function consider now the solution of equation in the form $u=e^{-i\alpha t}$ (65) Substituting equation (64) in equation (65) yield $-i^2 \hbar \alpha u = E u$ (66) Thus $E = \hbar \alpha$ (67) The periodicity condition requires $u(t + T) = u(t)$ (68) $e^{-i\alpha(t+T)} = e$ (69) $e^{-i\alpha T} = \cos \alpha T - i \sin \alpha T = 1$ (70) Thus $\cos \alpha T = 1$ $\sin \alpha T = 0$ (71) Hence $\alpha T = 2 n \pi$ (72) $\alpha = \frac{2n\pi}{\pi}$ $\frac{\partial f}{\partial T} = 2n\pi f = n\omega$ (73) In view of equations (66) and (67) the energy is thus given by $E = \hbar \alpha = n \hbar \omega$ (74) This means that the energy of gravity field is quantized. **ANGULAR MOMENTUM QUANTIZATION** Equation (23) can be used to find the relation $-\frac{c\hbar}{i\hbar}\frac{d\Phi}{c^2} = C_2$ (75)

ɸ − ℏ \mathcal{C} ɸ $\frac{d\Phi}{\varphi} = \frac{C_2}{c}$ \mathcal{C} (76)

[Abdallah* *et al.,* **6(4): April, 2017] Impact Factor: 4.116 IC[™] Value: 3.00 CODEN: IJESS7** $\hat{L}_z \Phi = L_z$ (77) Where the L_z operator is given by $\hat{L}_z = -\frac{\hbar}{i}$ i $\frac{\partial}{\partial \varphi}$ (78) $\int \frac{d\Phi}{\Gamma}$ $rac{d\Phi}{\Phi} = -\frac{i}{\hbar}$ ℏ \mathcal{C}_2 $\left(\frac{z}{c}\right) d\varphi$ (79) ln $\phi = -\frac{iC_2}{\sigma^2}$ $\frac{\partial^2 Z}{\partial t} = -iC_0 \varphi + C_3$ (80) $\Phi = e^{c_3} e^{-ic_0 \varphi} = \Phi_0 e^{-ic_0}$ (81) $\Phi(\varphi + 2\pi) = \Phi(\varphi)$ (82) $e^{-ic_0(\varphi + 2\pi)} = e^{ic_0\varphi}$ (83) $e^{-2\pi c_0 i} = \cos 2\pi c_0 + i \sin 2\pi c_0 = 1$ (84) Thus $\cos 2\pi c_0 = 1$ $\sin 2\pi c_0 = 0$ (85)
 $2\pi c_0 = 2n\pi$ (86) $2\pi C_0 = 2n\pi$ $n = 0, 1, 2, 3, \dots$... $C_0 = n$ $L_z = \frac{C_2}{c}$ $\frac{c_2}{c} = \frac{\hbar c}{c}$ $\frac{1}{c}C_0 = \hbar n$ (87)

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A full GGR quantum equation for spherically symmetric motion is obtained in equation (17). The separation of variables (18) leads to 4 independent (r, θ, ϕ, t) equations. The time dependent and radial parts [see equation (5)] predict. Again the existence of gravitational wave. Applying time periodicity of this wave on equation (15). Again this energy is nothing but quantum plank ordinary energy. The angular. ϕ Part of Schrödinger equation is given by equation (1). Its solution is given by equation (7). Using the fact that the wave function has only one unique value, the angular momentum component L_z is shown to be quantized.

CONCLUSION

The GGR quantum equation is useful in quantizing energy and angular momentum component along z direction. It predicts the existence of graviton and gravitational Waves.

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- [5] [Einstein 1917,](https://en.wikipedia.org/wiki/General_relativity#CITEREFEinstein1917) cf. [Pais 1982,](https://en.wikipedia.org/wiki/General_relativity#CITEREFPais1982) ch. 15e
- [6] Hubble's original article is [Hubble 1929;](https://en.wikipedia.org/wiki/General_relativity#CITEREFHubble1929) an accessible overview is given in [Singh 2004,](https://en.wikipedia.org/wiki/General_relativity#CITEREFSingh2004) ch. 2–4
- [7] As reported in [Gamow 1970.](https://en.wikipedia.org/wiki/General_relativity#CITEREFGamow1970) Einstein's condemnation would prove to be premature, cf. the section [Cosmology,](https://en.wikipedia.org/wiki/General_relativity#Cosmology) below
- [8] [Pais 1982,](https://en.wikipedia.org/wiki/General_relativity#CITEREFPais1982) pp. 253–25
- [9] [Kennefick 2005,](https://en.wikipedia.org/wiki/General_relativity#CITEREFKennefick2005) [Kennefick 2007](https://en.wikipedia.org/wiki/General_relativity#CITEREFKennefick2007)
- [10] Pais 1982, ch. 16
- [11]*Thorne, Kip (2003). [The future of theoretical physics and cosmology: celebrating Stephen Hawking's](https://books.google.com/books?id=yLy4b61rfPwC) [60th birthday.](https://books.google.com/books?id=yLy4b61rfPwC) Cambridge University Press. p. 74. [ISBN](https://en.wikipedia.org/wiki/International_Standard_Book_Number) [0-521-82081-2.](https://en.wikipedia.org/wiki/Special:BookSources/0-521-82081-2)*
- [12] Israel 1987, ch. 7.8–7.10, [Thorne 1994,](https://en.wikipedia.org/wiki/General_relativity#CITEREFThorne1994) ch. 3–
- [13] Sections [Orbital effects and the relativity of direction,](https://en.wikipedia.org/wiki/General_relativity#Orbital_effects_and_the_relativity_of_direction) [Gravitational time dilation and frequency shift](https://en.wikipedia.org/wiki/General_relativity#Gravitational_time_dilation_and_frequency_shift) and [Light deflection and gravitational time delay,](https://en.wikipedia.org/wiki/General_relativity#Light_deflection_and_gravitational_time_delay) and references therei
- [14] Section [Cosmology](https://en.wikipedia.org/wiki/General_relativity#Cosmology) and references therein; the historical development is in
- [15] The following exposition re-traces that of [Ehlers 1973,](https://en.wikipedia.org/wiki/General_relativity#CITEREFEhlers1973) sec.
- [16] Arnold 1989, ch. 1

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- [18] Will 1993, sec. 2.4[, Will 2006,](https://en.wikipedia.org/wiki/General_relativity#CITEREFWill2006) sec. 2
- [19] Wheeler 1990, ch.
- [20[\]Ehlers 1973,](https://en.wikipedia.org/wiki/General_relativity#CITEREFEhlers1973) sec. 1.2, [Havas 1964,](https://en.wikipedia.org/wiki/General_relativity#CITEREFHavas1964) [Künzle 1972.](https://en.wikipedia.org/wiki/General_relativity#CITEREFK.C3.BCnzle1972) The simple thought experiment in question was first described in

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- [21[\]Ehlers 1973,](https://en.wikipedia.org/wiki/General_relativity#CITEREFHeckmannSch.C3.BCcking1959) pp. 10f
- [22]Good introductions are, in order of increasing presupposed knowledge of mathematics, [Giulini 2005,](https://en.wikipedia.org/wiki/General_relativity#CITEREFGiulini2005) [Mermin 2005,](https://en.wikipedia.org/wiki/General_relativity#CITEREFMermin2005) an[d Rindler 1991;](https://en.wikipedia.org/wiki/General_relativity#CITEREFRindler1991) for accounts of precision experiments, cf. part IV of
- [23]An in-depth comparison between the two symmetry groups can be found in [Giulini 2006a](https://en.wikipedia.org/wiki/General_relativity#CITEREFGiulini2006a)
- [24[\]Rindler 1991,](https://en.wikipedia.org/wiki/General_relativity#CITEREFRindler1991) sec. 22, [Synge 1972,](https://en.wikipedia.org/wiki/General_relativity#CITEREFSynge1972) ch. 1 and 2
- [25[\]Ehlers 1973,](https://en.wikipedia.org/wiki/General_relativity#CITEREFEhlers1973) sec. 2.3
- [26[\]Ehlers 1973,](https://en.wikipedia.org/wiki/General_relativity#CITEREFEhlers1973) sec. 1.4, [Schutz 1985,](https://en.wikipedia.org/wiki/General_relativity#CITEREFSchutz1985) sec. 5.1
- [27[\]Ehlers 1973,](https://en.wikipedia.org/wiki/General_relativity#CITEREFEhlers1973) pp. 17ff; a derivation can be found in [Mermin 2005,](https://en.wikipedia.org/wiki/General_relativity#CITEREFMermin2005) ch. 12. For the experimental evidence, cf. the section [Gravitational time dilation and frequency shift,](https://en.wikipedia.org/wiki/General_relativity#Gravitational_time_dilation_and_frequency_shift) below
- [28] Rindler 2001, sec. 1.13; for an elementary account, see [Wheeler 1990,](https://en.wikipedia.org/wiki/General_relativity#CITEREFWheeler1990) ch. 2; there are, however, some differences between the modern version and Einstein's original concept used in the historical derivation of general relativity, cf. [Norton 1985](https://en.wikipedia.org/wiki/General_relativity#CITEREFNorton1985)
- [29[\]Ehlers 1973,](https://en.wikipedia.org/wiki/General_relativity#CITEREFEhlers1973) sec. 1.4 for the experimental evidence, see once more section [Gravitational time dilation](https://en.wikipedia.org/wiki/General_relativity#Gravitational_time_dilation_and_frequency_shift) [and frequency shift.](https://en.wikipedia.org/wiki/General_relativity#Gravitational_time_dilation_and_frequency_shift) Choosing a different connection with non-zero [torsion](https://en.wikipedia.org/wiki/Torsion_tensor) leads to a modified theory known as [Einstein–Cartan theory](https://en.wikipedia.org/wiki/Einstein%E2%80%93Cartan_theory)
- [30[\]Ehlers 1973,](https://en.wikipedia.org/wiki/General_relativity#CITEREFEhlers1973) p. 16, [Kenyon 1990,](https://en.wikipedia.org/wiki/General_relativity#CITEREFKenyon1990) sec. 7.2, [Weinberg 1972,](https://en.wikipedia.org/wiki/General_relativity#CITEREFWeinberg1972) sec. 2.8
- [31[\]Ehlers 1973,](https://en.wikipedia.org/wiki/General_relativity#CITEREFEhlers1973) pp. 19–22; for similar derivations, see sections 1 and 2 of ch. 7 in [Weinberg 1972.](https://en.wikipedia.org/wiki/General_relativity#CITEREFWeinberg1972) The Einstein tensor is the only divergence-free tensor that is a function of the metric coefficients, their first and second derivatives at most, and allows the spacetime of special relativity as a solution in the absence of sources of gravity, cf. [Lovelock 1972.](https://en.wikipedia.org/wiki/General_relativity#CITEREFLovelock1972) The tensors on both side are of second rank, that is, they can each be thought of as 4×4 matrices, each of which contains ten independent terms; hence, the above represents ten coupled equations. The fact that, as a consequence of geometric relations known a[s Bianchi](https://en.wikipedia.org/wiki/Bianchi_identities) [identities,](https://en.wikipedia.org/wiki/Bianchi_identities) the Einstein tensor satisfies a further four identities reduces these to six independent equations, e.g. [Schutz 1985,](https://en.wikipedia.org/wiki/General_relativity#CITEREFSchutz1985) sec. 8.3
- [32] Kenyon 1990, sec. 7.4
- [33] Brans & Dicke 1961, [Weinberg 1972,](https://en.wikipedia.org/wiki/General_relativity#CITEREFWeinberg1972) sec. 3 in ch. 7, [Goenner 2004,](https://en.wikipedia.org/wiki/General_relativity#CITEREFGoenner2004) sec. 7.2, and [Trautman 2006,](https://en.wikipedia.org/wiki/General_relativity#CITEREFTrautman2006) respectively
- [34] Wald 1984, ch. 4[, Weinberg 1972,](https://en.wikipedia.org/wiki/General_relativity#CITEREFWeinberg1972) ch. 7 or, in fact, any other textbook on general relativity
- [35] At least approximately, cf. [Poisson 2004](https://en.wikipedia.org/wiki/General_relativity#CITEREFPoisson2004)
- [36[\]Wheeler 1990,](https://en.wikipedia.org/wiki/General_relativity#CITEREFWheeler1990) p. xi
- [37[\]Wald 1984,](https://en.wikipedia.org/wiki/General_relativity#CITEREFWald1984) sec. 4.4
- [38] Wald 1984, sec. 4.1
- [39]For the (conceptual and historical) difficulties in defining a general principle of relativity and separating it from the notion of general covariance, see [Giulini 2006b](https://en.wikipedia.org/wiki/General_relativity#CITEREFGiulini2006b)
- [40] section 5 in ch. 12 of [Weinberg 1972](https://en.wikipedia.org/wiki/General_relativity#CITEREFWeinberg1972)
- [41]Introductory chapters of [Stephani et al. 2003](https://en.wikipedia.org/wiki/General_relativity#CITEREFStephaniKramerMacCallumHoenselaers2003)
- [42]A review showing Einstein's equation in the broader context of other PDEs with physical significance is [Geroch 1996](https://en.wikipedia.org/wiki/General_relativity#CITEREFGeroch1996)
- [43]For background information and a list of solutions, cf. [Stephani et al. 2003;](https://en.wikipedia.org/wiki/General_relativity#CITEREFStephaniKramerMacCallumHoenselaers2003) a more recent review can be found in [MacCallum 2006](https://en.wikipedia.org/wiki/General_relativity#CITEREFMacCallum2006)
- [44[\]Chandrasekhar 1983,](https://en.wikipedia.org/wiki/General_relativity#CITEREFChandrasekhar1983) ch. 3,5,6
- [45[\]Narlikar 1993,](https://en.wikipedia.org/wiki/General_relativity#CITEREFNarlikar1993) ch. 4, sec. 3.3